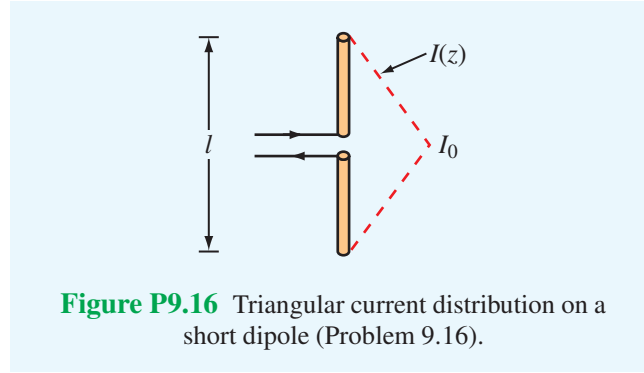


**9.16** For a short dipole with length  $l$  such that  $l \ll \lambda$ , instead of treating the current  $\tilde{I}(z)$  as constant along the dipole, as was done in Section 9-1, a more realistic approximation that ensures the current goes to zero at the dipole ends is to describe  $\tilde{I}(z)$  by the triangular function

$$\tilde{I}(z) = \begin{cases} I_0(1 - 2z/l), & \text{for } 0 \leq z \leq l/2 \\ I_0(1 + 2z/l), & \text{for } -l/2 \leq z \leq 0 \end{cases}$$

as shown in Fig. P9.14. Use this current distribution to determine the following:

- (a) The far-field  $\tilde{\mathbf{E}}(R, \theta, \phi)$ .
- (b) The power density  $S(R, \theta, \phi)$ .
- (c) The directivity  $D$ .
- (d) The radiation resistance  $R_{\text{rad}}$ .



**Figure P9.16** Triangular current distribution on a short dipole (Problem 9.16).

**Solution:**

(a) When the current along the dipole was assumed to be constant and equal to  $I_0$ , the vector potential was given by Eq. (9.3) as:

$$\tilde{\mathbf{A}}(R) = \hat{\mathbf{z}} \frac{\mu_0}{4\pi} \left( \frac{e^{-jkR}}{R} \right) \int_{-l/2}^{l/2} I_0 dz.$$

If the triangular current function is assumed instead, then  $I_0$  in the above expression should be replaced with the given expression. Hence,

$$\tilde{\mathbf{A}} = \hat{\mathbf{z}} \frac{\mu_0}{4\pi} \left( \frac{e^{-jkR}}{R} \right) I_0 \left[ \int_0^{l/2} \left( 1 - \frac{2z}{l} \right) dz + \int_{-l/2}^0 \left( 1 + \frac{2z}{l} \right) dz \right] = \hat{\mathbf{z}} \frac{\mu_0 I_0 l}{8\pi} \left( \frac{e^{-jkR}}{R} \right),$$

which is half that obtained for the constant-current case given by Eq. (9.3). Hence, the expression given by (9.9a) need only be modified by the factor of 1/2:

$$\tilde{\mathbf{E}} = \hat{\boldsymbol{\theta}} \tilde{E}_\theta = \hat{\boldsymbol{\theta}} \frac{jI_0 l k \eta_0}{8\pi} \left( \frac{e^{-jkR}}{R} \right) \sin \theta.$$

(b) The corresponding power density is

$$S(R, \theta) = \frac{|\tilde{E}_\theta|^2}{2\eta_0} = \left( \frac{\eta_0 k^2 I_0^2 l^2}{128\pi^2 R^2} \right) \sin^2 \theta.$$

(c) The power density is 4 times smaller than that for the constant current case, but the reduction is true for all directions. Hence,  $D$  remains unchanged at 1.5.

(d) Since  $S(R, \theta)$  is 4 times smaller, the total radiated power  $P_{\text{rad}}$  is 4-times smaller. Consequently,  $R_{\text{rad}} = 2P_{\text{rad}}/I_0^2$  is 4 times smaller than the expression given by Eq. (9.35); that is,

$$R_{\text{rad}} = 20\pi^2 (l/\lambda)^2 \quad (\Omega).$$

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